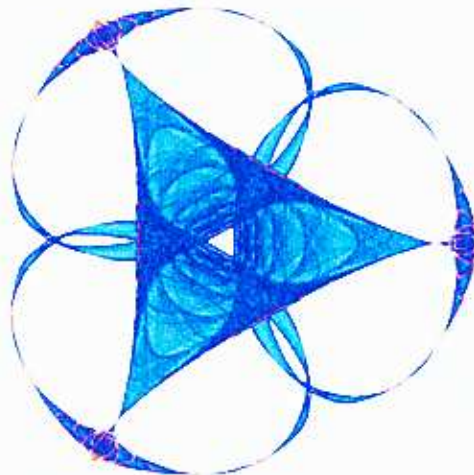


FUNCTORIAL REMARKS ON THE GENERAL CONCEPT OF CHAOS

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F.W. LAWVERE



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**INSTITUTE FOR MATHEMATICS AND ITS APPLICATIONS**

**UNIVERSITY OF MINNESOTA**

**514 Vincent Hall**

**206 Church Street S.E.**

**Minneapolis, Minnesota 55455**

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By

F. William Lawvere  
Department of Mathematics  
SUNY at Buffalo  
Buffalo, New York 14222

# Functorial Remarks on the General Concept of Chaos

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F. W. Lawvere

Although the term "chaos" is employed in various ways in current dynamics literature, several instances [1], [2] of a precise usage have in common two features of a categorical nature: surjectivity of an induced map and right adjointness in the inducing process. Abstracting these features we propose a general definition and then note that there are examples which are rather different from the usual ones but which nonetheless have an intuitively "chaotic" quality.

Recall that a functor

$$\mathcal{X} \xrightarrow{U} \mathcal{Y}$$

between two categories is said to have a right adjoint

$$\mathcal{Y} \xrightarrow{H} \mathcal{X}$$

if there are natural transformations  $1_{\mathcal{X}} \xrightarrow{\eta} H U$ ,

$U H \xrightarrow{\theta} 1_{\mathcal{Y}}$  which induce bijections

$$\frac{U X \longrightarrow Y}{X \longrightarrow H Y}$$

between the indicated sets of  $\mathcal{Y}$ -morphisms and  $\mathcal{X}$ -morphisms,

for each  $X$  in  $\mathcal{X}$  and  $Y \in \mathcal{Y}$ . An important class of examples may be constructed as follows. Let  $\mathcal{Y}$  be a suitable category of topological or differentiable spaces, and let  $T$  be a monoid in  $\mathcal{Y}$ . For example  $T$  could be the additive monoid of nonnegative reals or of nonnegative integers. Define  $\mathcal{X} = \mathcal{Y}^T$ , the category whose objects are objects of  $\mathcal{Y}$  equipped with actions of  $T$ , and whose morphisms are  $T$ -equivariant maps  $\varphi$ :

$$\begin{array}{ccc}
 \underline{X} \times T & \xrightarrow{\varphi \times T} & Y \times T \\
 \text{action} \downarrow & & \downarrow \text{action} \\
 \underline{X} & \xrightarrow{\varphi} & Y
 \end{array}
 \quad \begin{array}{l} \text{commutative} \\ \text{in } \mathcal{Y} \end{array}$$

We will write the actions of  $T$  multiplicatively and on the right. Define  $U$  to be the "underlying space" functor which forgets the action

$$\begin{array}{ccc}
 \mathcal{X} & \xrightarrow{U} & \mathcal{Y} \\
 \parallel & \nearrow \text{forget} & \\
 \mathcal{Y}^T & & 
 \end{array}$$

As usual we will omit the symbol  $U$  from the notation when it causes no confusion. Then  $U$  has a right adjoint given by,

$$H(Y) = Y^T$$

the internal function space (presumed to exist in  $\mathcal{Y}$ ) for each space  $Y$  in  $\mathcal{Y}$ . The adjointness is demonstrated

in essence by the bijection

$$\frac{X \xrightarrow{\bar{\varphi}} Y^T}{X \xrightarrow{\varphi} Y}$$

between arbitrary maps  $\varphi$  and equivariant maps  $\bar{\varphi}$  defined by

$$\varphi(x) = \bar{\varphi}(x)(u)$$

where  $u$  is the unit element of  $T$

$$\bar{\varphi}(x)(t) = \varphi(xt)$$

where  $xt$  is the action given on  $X$

Here  $\bar{\varphi}$  is equivariant if  $\varphi$  is given and we have equipped  $Y^T$  with the usual action by "translation":

$$(yt)(s) = y(ts) \quad \text{for all } s \in T, y \in Y^T$$

Thus if  $X$  is a space equipped with a dynamical action of  $T$  and  $Y$  is a given space (without action) and if we consider a map  $X \xrightarrow{\varphi} Y$  as an "observable" of state, then  $\bar{\varphi}(x)$  is the function of time giving the progression of observed values of  $\varphi$  if we start in state  $x$  at  $u =$  time  $0$ . This  $\bar{\varphi}$  is sometimes referred to as "symbolic dynamics", the points of  $Y$  being considered as symbols for the blocks into which  $\varphi$  divides the state space  $X$ ; then starting in state  $x$ ,  $\bar{\varphi}(x)$  is the  $T$ -sequence of blocks through which the dynamics takes the system.

Definition If  $X \in Y^T$ ,  $Y \in Y$ , an observable  $X \xrightarrow{\varphi} Y$  is  $T$ -chaotic iff the induced map  $X \xrightarrow{\bar{\varphi}} Y^T$  is an epimorphism. More generally if  $\mathcal{X} \xrightarrow{v} \mathcal{Y}$  is a functor with right adjoint  $H$  and  $UX \xrightarrow{\varphi} Y$ ,  $\varphi$  is  $U$ -chaotic iff  $X \xrightarrow{\bar{\varphi}} H(Y)$  is an epimorphism.

That is (for those categories  $\mathcal{Y}$  in which epimorphisms are surjective)  $X \xrightarrow{\varphi} Y$  is a chaotic observable iff every  $T$ -sequence of symbols  $T \xrightarrow{y} Y$  is realized as  $y = \bar{\varphi}(x)$  for at least one state  $x \in \bar{X}$ . Refinements of this condition can be formulated by considering certain homomorphisms  $T \rightarrow T'$  as periods and requiring the same condition after applying the corresponding change of action functors.

A standard example of a map  $\varphi$  often considered in this context is the adjunction

$$X \longrightarrow \hat{\pi}_0(X) = Y$$

where  $\hat{\pi}_0$  is the "space of components" functor left adjoint to the inclusion  $\mathcal{Y}_0 \hookrightarrow \mathcal{Y}$  of the category of prodiscrete spaces. In case  $Y = \hat{\pi}_0(X)$  is actually finite and  $T$  is countable discrete, then  $H(Y) = Y^T$  is a Cantor space.

Finally, in accord with loose usage we may say that a system  $\bar{X}$  in  $\mathcal{X} = Y^T$  is "chaotic" in case there exists an  $\mathcal{X}$ -subobject  $X \hookrightarrow \bar{X}$  (that is, a  $T$ -invariant subspace) and a nontrivial space  $Y$  and a chaotic  $X \rightarrow Y$ , or perhaps that  $\hat{\pi}_0(X)$  is nontrivial and  $X \rightarrow \hat{\pi}_0 X$  chaotic.

The extent to which the above general categorical definition expresses the informal notion of chaos can be illustrated by an example of a dual algebraic nature. Let  $\mathcal{Y} = \mathcal{a}$  = the category of all commutative algebras over the reals; for example the function ring  $C^\infty(Y)$  is an object of  $\mathcal{a}$  for a manifold  $Y$ . Let  $\mathcal{X} = \mathcal{a}'$  = the category of all such algebras  $A$  equipped moreover with derivations  $AD(\cdot)$ ; that is  $(\cdot)'$  is a real-linear map satisfying the Leibniz product rule, where the morphisms of  $\mathcal{a}'$  are algebra homomorphisms which moreover commute with the given derivations. Of course any vector field supplies an example of an object in  $\mathcal{a}'$ . Then the forgetful functor

$$\mathcal{a}' \xrightarrow{U} \mathcal{a}$$

has the right adjoint  $H$  which assigns, to any algebra  $A$ , the algebra

$$H(A) = A[[t]]$$

of all formal (divided) power series with coefficients in  $A$ , equipped with the obvious formal derivation  $\frac{d}{dt}$ . The adjointness is verified as follows: If  $B$  is equipped with any given derivation and  $B \xrightarrow{g} A$  is an algebra homomorphism, then

$$B \xrightarrow{g} A[[t]]$$

is defined by  $\bar{\varphi}(f) = \sum \frac{\varphi(f^{(n)})}{n!} t^n$   
 where  $f^{(n)}$  is defined by iterating the  
 given derivation on  $B$ , and clearly  $\bar{\varphi}(f') = \frac{d}{dt} \bar{\varphi}(f)$   
 as required for  $a'$ -morphisms.

Proposition: If  $B = C^\infty(\Sigma)$ , where  $\Sigma$  is a  
 manifold, is equipped with the derivation induced by any  
 given non-trivial vector field, then  $B$  has a map

$$B \xrightarrow{\varphi} \mathbb{R} \text{ which is chaotic relative to the adjunction } a' \hookrightarrow a.$$

Proof: Let  $X$  be a point where the vector field is non-  
 trivial and let  $\varphi(f) = f(X)$  for all  $f \in B$ . Then  $\bar{\varphi}(f)$   
 is the Taylor series of  $f$ . Borel's theorem states that  
any element of  $H(\mathbb{R}) = \mathbb{R}[[t]]$  is the Taylor series at  $X$  of  
 some smooth  $f$ . Thus  $\bar{\varphi}$  is surjective, so that  $\varphi$  is  
 chaotic relative to  $a' \rightarrow a$ , according to our definition.

Some other simple examples of right adjoints which  
 may be of interest in this connection are determined by  
 starting with a homomorphism  $T' \rightarrow T$  of monoids in a  
 special category and considering the induced-action  
 functor

$$S^T \xrightarrow{U} S^{T'}$$

whose right adjoint (generalizing ( ) when =1) is

$$H(Y) = \text{Hom}_{T'}(T, Y)$$



the space of all  $T'$ -equivariant maps  $T \rightarrow Y$ , where  $Y$  has some given  $T'$ -action and  $T'$  acts on  $T$  via translation and the given  $T' \rightarrow T$ . For example if we consider the inclusion  $\mathbb{N} \hookrightarrow \mathbb{R}^+$  of additive monoids,  $Y$  is essentially equipped with a single  $\mathcal{S}$ -endomorphism  $\tau$  (the action of  $1 \in \mathbb{N}$ ) and so is a "discrete time dynamical system", and the associated continuous-time dynamical system  $H(Y)$  has as its states all  $\mathcal{S}$ -maps  $\mathbb{R}^+ \rightarrow Y$  for which

$$f(t+1) = f(t) \cdot \tau \quad \text{for all } t \geq 0$$

A morphism  $\bar{X} \xrightarrow{\varphi} Y$ , where  $\bar{X} \in \mathcal{S}^{\mathbb{R}^+}$  is now required to satisfy the condition

$$\varphi(X \cdot t) = \varphi(X) \tau^t$$

in case  $t=n$  is a whole number. [We remark that, as in many examples, our  $\mathcal{U}$  has also a left adjoint  $Y \mapsto \mathbb{R}^+_{\mathbb{N}} Y$ , a tensor-like quotient space of  $\mathbb{R} \times Y$ , which gives a different notion of "continuous-time system associated to a discrete-time system" ] Such a morphism  $\varphi$  is chaotic (in the present relative sense) if and only if every  $f$  of the kind described above are, for some  $X \in \bar{X}$  of the form

$$f(t) = \varphi(X \cdot t) \quad \text{for all } t.$$

Of course there are standard examples of such  $\varphi$ , namely the evaluation at  $0: \bar{X} = H(Y) \rightarrow Y$ , but the novelty would be to find  $\bar{X}$  constructed by other means.

In the case of a surjective  $T' \rightarrow T$ , such as  
 $\mathbb{R} \rightarrow \mathbb{R}/\mathbb{Z}$  induced by a given period,  
 $H(Y) \hookrightarrow Y$

is simply the subspace of the  $T'$ -system  $Y$  consisting of points having that period, so for  $X \xrightarrow{\varphi} Y$  to be "chaotic" relative to  $T' \rightarrow T$  simply means that all those points of  $Y$  having the period in question are values of  $\varphi$ .

As a final example, consider the problem of solving differential equations. From the point of view of "Synthetic Differential Geometry" [3] this can be considered as a special case of the special class of adjoints arising from a change of monoid  $\langle D \rangle \rightarrow \mathbb{R}$ , but in any case, if  $\mathcal{S}$  is a reasonable category of differentiable spaces, we can consider the category  $\mathcal{Y}$  whose typical object  $Y$  is a space in  $\mathcal{S}$  equipped with a vector field, and whose morphisms are smooth maps  $Y_1 \rightarrow Y_2$  which commute with the designated vector fields. Then if  $\mathcal{X} = \mathcal{S}^{\mathbb{R}}$  is the category of reversible continuous-time flows, there is a functor  $\mathcal{X} \xrightarrow{\mathcal{U}} \mathcal{Y}$  given by differentiating each flow at  $t=0$ . The right adjoint to  $\mathcal{U}$  assigns to each such differential equation  $Y$  the space  $H(Y)$  of all solution curves which are defined for all time. Thus for the adjunction map  $H(Y) \rightarrow Y$  to be "chaotic" in this context merely means that through each

point of  $Y$  there is a solution curve which extends for all time. More generally if  $\bar{X}$  is a flow and  $Y$  a vector field, a compatible map  $X \xrightarrow{\mathcal{P}} Y$  is "chaotic" relative to the differentiation functor  $U$  iff every solution curve for  $Y$  is for some  $x$  the  $\mathcal{P}$ -image of the flow through  $x$ .

That the last few examples sound more like normal, reasonable behavior than like pathological chaos is explained in terms of the extent to which the functor  $U$ , or more particularly the monoid homomorphism  $T' \rightarrow T$ , is an isomorphism: In the basic example  $1 \rightarrow T, Y \mapsto H(Y) = Y^T$  is a big jump, so surjectivity of  $\bar{\mathcal{P}}$  is harder to come by. In the case of  $\langle D \rangle \hookrightarrow \mathbb{R}$ , the functor  $U$  is not quite an equivalence but reasonable  $Y$  "think" it is in that  $H(Y) \xrightarrow{\sim} Y$  due to the strong existence and uniqueness property; thus for such  $Y$ , the surjectivity of  $\mathcal{P}$  itself is sufficient for  $\mathcal{P}$  to be "chaotic".

- [1] Saari, D.G., and Urenko, J.B., "Newton's method, Circle Maps, and Chaotic Motion", American Mathematical Monthly, Jan. 1984
- [2] Coppel, W.A., "Maps on an Interval", IMA Preprint Series # 26, June 1983, Institute for Mathematics and its Applications, University of Minnesota.
- [3] Kock, A. Synthetic Differential Geometry, LMS Lecture Note Series # 51, Cambridge University Press 1981.